A Remarkable Elevation of Freezing Temperature of CCl₄ in Graphitic Micropores

Katsumi Kaneko,* Ayumi Watanabe, Taku Iiyama, Ravi Radhakrishnan,† and Keith E. Gubbins†

Physical Chemistry, Material Science, Graduate School of Science and Technology, Chiba University, 1-33 Yayoi, Inage, Chiba 263-8522, Japan and Chemical Engineering, North Carolina State University, 113 Riddick Labs, Raleigh, North Carolina 27695-7905

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The freezing behavior of CCl₄ confined in graphitic micropores of activated carbon fibers (ACFs) was examined by use of differential scanning calorimetry over the temperature range of 180 to 320 K. For ACF of average pore width, \( w \) = 1.1 nm, we observed a clear peak showing that the freezing temperature is 299 K. The freezing temperature was elevated by 57 K. However, the enthalpy of freezing was only 1.0% of that of the bulk liquid CCl₄ to an fcc solid phase. CCl₄ confined in micropores of \( w \) = 0.80 nm had a more diffuse DSC peak, providing a similar elevation of the freezing temperature and an even smaller freezing enthalpy of 0.5% of the bulk value.

Introduction

The structure and phase behavior of molecular assemblies confined in small pores have attracted much attention. Freezing and melting of condensates in mesopores have been studied for a long time.1 Recently, well-defined porous materials have led to an active study on the phase behavior of confined fluids.2–5 In almost all cases a clear depression of the freezing temperature was observed, and the depression has been basically explained by the capillary effect in surface chemistry. Recent experimental and molecular simulation studies have shown that the classical capillary theory is not enough.6–8 Miyahara and Gubbins8 simulated freezing and melting phenomena for a Lennard–Jones model of CH₄ in graphitic slit pores. They showed that the freezing temperature is elevated relative to that of the bulk liquid; the increase is greater for smaller pores, but the freezing disappears in small micropores.7,8 This predicted elevation of freezing in porous carbons has recently been reconfirmed by careful free energy calculations by Radhakrishnan and Gubbins.9 However, we do not have sufficient experimental data on the phase behavior of fluids confined in micropores for understanding this phenomenon. Organized structures have been reported for CCl₄,10,11 H₂O,12 SO₂,13 and O₂14 in micrographitic slit micropores of activated carbon fiber (ACF), which change with the slit width. Hence, the micropores of ACFs can offer a good system for studying the freezing effect in confined structures.

The structures of CCl₄ in ACF micropores at 303 K were examined by in situ X-ray diffraction and GCMC simulation, showing that the CCl₄ molecular assembly structure depends on the pore width.10,11 Also, the freezing of bulk CCl₄ had been studied using a calorimetric method and X-ray diffraction.15–18 Information on bulk CCl₄ should be helpful for understanding CCl₄ in micropores.

This letter reports a marked elevation of the freezing temperature of CCl₄ in micrographitic micropores.

Experimental and Simulation Section

Pitch-based ACFs of different pore widths (P5, P10, and P20) were used as the micrographitic micropore system. The mi-
TABLE 1: Micropore Structure of ACF Samples by N2 Adsorption

<table>
<thead>
<tr>
<th>micropore vol (mL g⁻¹)</th>
<th>surface area (m² g⁻¹)</th>
<th>pore width (nm)</th>
</tr>
</thead>
<tbody>
<tr>
<td>P5</td>
<td>0.336</td>
<td>0.75</td>
</tr>
<tr>
<td>P10</td>
<td>0.622</td>
<td>0.82</td>
</tr>
<tr>
<td>P20</td>
<td>0.970</td>
<td>1.13</td>
</tr>
</tbody>
</table>

Figure 1. DSC cooling curves of CCl4 confined in micropores of various ACFs and of bulk CCl4 liquid: (a) bulk liquid, (b) P20, (c) P10, and (d) P5. Here the ordinate scale of bulk liquid is reduced by 1/10.

cropore volume, the specific surface area, and the slit pore width were determined by the high-resolution N2 adsorption isotherm at 77 K using a subtracting pore effect method.19,20 Table 1 lists the micropore structural parameters. The pore widths were in the range of 0.75 to 1.1 nm. CCl4 was adsorbed at 303 K under a saturated vapor pressure after drying ACF samples by evacuation at 373 K for 2 h. The ACF samples with adsorbed CCl4 were sealed in an aluminum pan. The DSC chart was recorded at the temperature scanning rate of 5 K min⁻¹ or 1 K min⁻¹ using a DSC system (MAC Science, DSC3100). We repeated measurements of the DSC cooling or heating charts three times and summed them to obtain the average DSC curve. The freezing temperature was determined by the crossing point of the linear rising part of the peak and the average background. The enthalpy of freezing was obtained from the peak area. The heat evolution peaks which are ascribed to the liquid to fcc solid phase (Ia) transition, the Ib to rhombohedral phase (Ib) transition, and the Ic to monoclinic phase (solidm) transition begin at 241.8 K, 228.4 K, and 218.3 K, respectively. These transition temperatures in the literature are 242 K, 231 K, and 225 K for the liquid–Ia, Ia–Ib, and Ib–solidm transitions, respectively.16–18

The observed transition temperatures are close to the literature values, although the correct determination of the phase transition temperature is difficult from the cooling curve because of supercooling. The peak of the Ib–solidm transition is the most intense. On the other hand, only one diffuse peak of the heat evolution for CCl4 confined in micropores of ACF is observed around 300 K, which is much higher than that of the bulk liquid. This peak should come from the freezing of CCl4 confined in micropores, although three peaks are not observed, as in the bulk CCl4. The strong confinement of CCl4 molecules in the micropores should prohibit ordinary solid-phase transitions. The DSC peaks for the P20 micropore material is more evident than those for the P10 and P5 materials. The freezing temperature of CCl4 for P20 is 299 K, being higher than that of bulk CCl4 by 57 K. Also, the freezing temperatures for P10 and P5 are estimated to be 301 and 300 K, respectively, from the quite diffuse noisy peak; in particular, the determination of the freezing temperature of P5 is quite difficult owing to strong diffuseness. The enthalpy of freezing, ΔHf, for the confined CCl4 is much smaller than that of the bulk liquid, and the ΔHm values are listed in Table 2. The ΔHf of the confined CCl4 becomes smaller with decreasing pore width, while the freezing temperature does not vary significantly with the pore width.

Figure 2 shows the DSC curves of CCl4 confined in micropores of P20 upon cooling (a) and heating (b). The heating curve of bulk solid CCl4 is shown for comparison, which has two peaks at 223 and 248 K originating from the solidm–Ib and Ib–liquid transitions.15 Figure 2 shows both the remarkable elevation of the melting temperature of confined CCl4 by 50 K and the markedly small enthalpy of melting. There is no meaningful difference between the freezing and melting temperatures of confined CCl4. The enthalpy of melting of confined CCl4 is quite small compared with the bulk value, as is the enthalpy of freezing. Thus, the melting temperature of confined CCl4 is much higher than that of bulk CCl4 liquid. DSC measurements upon heating were carried out on other ACF samples; similar tendencies were observed, as shown in Table 2. The melting temperature of confined CCl4 does not depend on the pore width, but the enthalpy of melting becomes smaller with decreasing the pore width in a similar way to the freezing behavior.

The features of the observed DSC results on CCl4 confined in graphitic micropores of which micropore width is less than 1.2 nm are as follows. The elevation of the freezing temperature

<table>
<thead>
<tr>
<th>adsorbedCCl4 (mg g⁻¹)</th>
<th>freezing temp (K)</th>
<th>ΔHf (J mol⁻¹)</th>
<th>melting temp (K)</th>
<th>ΔHm (J mol⁻¹)</th>
</tr>
</thead>
<tbody>
<tr>
<td>P5</td>
<td>330</td>
<td>300</td>
<td>5</td>
<td>299</td>
</tr>
<tr>
<td>P10</td>
<td>555</td>
<td>301</td>
<td>8</td>
<td>300</td>
</tr>
<tr>
<td>P20</td>
<td>1470</td>
<td>299</td>
<td>17</td>
<td>298</td>
</tr>
<tr>
<td>bulk liquid</td>
<td>242</td>
<td>1570</td>
<td></td>
<td></td>
</tr>
<tr>
<td>bulk solid</td>
<td>248</td>
<td>2190</td>
<td></td>
<td></td>
</tr>
</tbody>
</table>

* Here, ΔHf and ΔHm denote the enthalpies of freezing and melting, respectively. The enthalpy of melting of bulk CCl4 in the literature is 2.56 kJ mol⁻¹.
CCl$_4$ is 0.5881 nm, and the intermolecular interaction is limited to the nearest neighbors. Hence, the interaction of a CCl$_4$ solid even at 303 K on the basis of X-ray diffraction data, in micropores has an ordered structure similar to that of the is formed in case of the pores of ACFs at 303 K from the X-ray diffraction patterns, showing that the second-nearest neighbor confined in micropores of ACFs at 303 K. The Lennard-Jones size parameter of CCl$_4$ molecules shifts to a smaller side because of the weaker than the interaction of a CCl$_4$ molecule with the graphitic pore wall. Therefore, the observed elevation of the freezing and melting temperatures for confined CCl$_4$ can be explained basically by the Miyahara and Gubbins theory. Iiyama et al. determined the electron radial distribution function of CCl$_4$ confined in micropores of ACFs at 303 K from the X-ray diffraction patterns, showing that the second-nearest neighbor peak of CCl$_4$ molecules shifts to a smaller side because of the more dense structure of confined CCl$_4$ and that the plastic crystalline structure, which is found at 253 K in the bulk phase, is formed in case of the pores of $w = 1.1$ nm (corresponding to P20). The upward shift obtained by the simulation is greater than the observed value. However, the agreement is reasonable regardless of the simple model. The simulation suggested the phase transition of the crystal phase to the orientationally ordered fluid upon freezing for confined CCl$_4$ in micropores of ACF, being consistent with the observed small enthalpy.

So far we have no direct structural information on the phase transition across the observed transition temperature using X-ray diffraction measurements. We hope to provide such data in the future after improvement of our X-ray diffraction system.

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References and Notes

Figure 2. Averaged DSC curves of CCl$_4$ confined in micropores of P20 on cooling and heating. The heating curve of bulk CCl$_4$ solid, whose ordinate is reduced to 1/10, is shown: (a) cooling and (b) heating curves of confined CCl$_4$, (c) heating curve of bulk CCl$_4$. is about 58 (±1) K, which does not sensitively depend on the pore width for this range of pore widths. The enthalpy of freezing is less than 1% of the bulk value. The melting of confined CCl$_4$ has a similar tendency to the freezing behavior. A confined liquid in mesopores should exhibit a depression of the freezing and melting temperatures according to classical surface chemistry. The phase behavior of confined CCl$_4$ in carbon is quite unique, being different from those of liquid confined in most other materials. Miyahara and Gubbins proposed that if the interaction of a molecule with the sample pore surface is greater than that with the pore wall consisting of adsorbate molecules, the freezing temperature should be elevated with decreasing pore width and the freezing disappears in narrow micropores. The Lennard-Jones size parameter of CCl$_4$ is 0.5881 nm, and the intermolecular interaction is limited to the nearest neighbors. Hence, the interaction of a CCl$_4$ molecule with the wall made of CCl$_4$ molecules should be weaker than the interaction of a CCl$_4$ molecule with the graphitic pore wall. Therefore, the observed elevation of the freezing and melting temperatures for confined CCl$_4$ can be explained basically by the Miyahara and Gubbins theory. Iiyama et al. determined the electron radial distribution function of CCl$_4$ confined in micropores of ACFs at 303 K from the X-ray diffraction patterns, showing that the second-nearest neighbor peak of CCl$_4$ molecules shifts to a smaller side because of the more dense structure of confined CCl$_4$ and that the plastic crystalline structure, which is found at 253 K in the bulk phase, is formed in case of the pores of $w = 0.75$ nm. Thus, CCl$_4$ in micropores has an ordered structure similar to that of the solid even at 303 K on the basis of X-ray diffraction data, coinciding with the observed DSC behavior. Hence, the freezing or melting temperature should be higher than the bulk value, as shown in this work. However, CCl$_4$ molecules in micropores cannot form a highly crystalline structure because of spatial restrictions. Also, even the liquid structure of confined CCl$_4$ should be more ordered than that of the bulk liquid. Thus the freezing of confined CCl$_4$ must be the transition of a partially ordered liquid to a highly defective solid; the enthalpy of freezing or melting in such a case is much smaller than the bulk value, as observed.

The simulation of the freezing of CCl$_4$ in a graphite slit showed that the freezing temperature is elevated by 90 K for pores of $w = 1.1$ nm (corresponding to P20). The upward shift obtained by the simulation is greater than the observed value. However, the agreement is reasonable regardless of the simple model. The simulation suggested the phase transition of the crystal phase to the orientationally ordered fluid upon freezing for confined CCl$_4$ in micropores of ACF, being consistent with the observed small enthalpy.

So far we have no direct structural information on the phase transition across the observed transition temperature using X-ray diffraction measurements. We hope to provide such data in the future after improvement of our X-ray diffraction system.