

Fourier transforms

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- ▶ Fourier analysis of discrete signals $x : [0, N-1] \rightarrow \mathbb{C} \Rightarrow \mathsf{DFT}$, iDFT
- Good (and quick) computational tool
 - \Rightarrow Signal analysis $\ \Rightarrow$ pattern discovery, frequency components
 - \Rightarrow Signal processing \Rightarrow compression, noise removal
- Two important limitations
 - \Rightarrow Time is neither discrete nor finite (not always, at least)
 - \Rightarrow Properties and interpretations are easier in continuous time
- Fourier analysis of continuous signals \Rightarrow Fourier transform (FT)



Continuous time signals

Fourier transform

Inverse Fourier transform

Delta function

Generalized orthogonality

Generalized Fourier transforms

Properties of the Fourier transform

Convolution



- ▶ We have been dealing with discrete signals $x : [0, N-1] \rightarrow \mathbb{C}$
- To infinity \Rightarrow Let number of samples go to infinity
 - \Rightarrow Discrete time signal $x : \mathbb{Z} \to \mathbb{C}$
 - \Rightarrow Values x(n) for $n = \ldots, -1, 0, 1, \ldots$
- ► And beyond ⇒ Fill in the gaps between samples
 - \Rightarrow Continuous time signal $x : \mathbb{R} \to \mathbb{C}$
 - \Rightarrow Values x(t) for t any real number in $(-\infty, +\infty)$
- Let's begin by studying continuous time signals



- Continuous time variable $t \in \mathbb{R}$.
- Continuous time signal x is a function that maps t to real value x(t)

$$x:\mathbb{R}\to\mathbb{R}$$

- The values that the signal takes at time t is x(t)
- It will make sense to talk about complex signals (as in discrete case)

$$x:\mathbb{R}\to\mathbb{C}$$

• where the values $x(t) = x_R(t) + j x_I(t)$ are complex numbers





► Given two signals x and y define the inner product of x and y as

$$\langle x,y \rangle := \int_{-\infty}^{\infty} x(t) y^*(t) dt$$

• Akin to inner product of discrete signals $\Rightarrow \langle x, y \rangle = \sum_{n=0}^{N} x(n)y(n)$



- But we have infinite number of components. To infinity and beyond
- Intuition holds $\Rightarrow \langle x, y \rangle$ is how much of y falls in x direction
- E.g., if $\langle x, y \rangle = 0$ the signals are orthogonal. They are "unrelated"



▶ As for regular (finite dimensional) signals define the norm of signal x

$$\|x\| := \left[\int_{-\infty}^{\infty} |x(t)|^2 dt\right]^{1/2} = \left[\int_{-\infty}^{\infty} |x_R(t)|^2 dt + \int_{-\infty}^{\infty} |x_I(t)|^2 dt\right]^{1/2}$$

▶ More important, define the energy of the signal as the norm squared

$$\|x\|^2 := \int_{-\infty}^{\infty} |x(t)|^2 dt = \int_{-\infty}^{\infty} |x_R(t)|^2 dt + \int_{-\infty}^{\infty} |x_I(t)|^2 dt$$

- For complex numbers $x(t)x^*(t) = |x_R(t)|^2 + |x_I(t)|^2 = |x(t)|^2$
- ▶ Thus, we can write the energy as $\Rightarrow \|x\|^2 = \langle x, x \rangle$
- Energy might be infinite. When energy is finite we write $||x||^2 < \infty$



> The largest an inner product can be is when the vectors are collinear

$$-\|x\| \|y\| \le \langle x, y \rangle \le \|x\| \|y\|$$

• Or in terms of energy $\Rightarrow \langle x, y \rangle^2 \le ||x||^2 ||y||^2$

If you are the sort of person that prefers explicit expressions

$$\int_{-\infty}^{\infty} x(t) y^*(t) dt \leq igg[\int_{-\infty}^{\infty} |x(t)|^2 dtigg] igg[\int_{-\infty}^{\infty} |y(t)|^2 dtigg]$$

▶ The equalities hold if and only if *x* and *y* are collinear



• The square pulse is the signal $\sqcap_T(t)$ that takes values

$$\Box_{T}(t) = 1 \qquad \text{for } -\frac{T}{2} \le t < \frac{T}{2}$$
$$\Box_{T}(t) = 0 \qquad \text{otherwise}$$
$$-\frac{T/2}{T/2} \qquad \frac{1}{T/2}$$

▶ To compute energy of the pulse we just evaluate the definition

$$\| \sqcap_{\mathcal{T}} (t) \|^2 := \int_{-\infty}^{\infty} |\sqcap_{\mathcal{T}} (t)(t)|^2 dt = \int_{-T/2}^{T/2} |1|^2 dt = \mathcal{T}$$

- Energy proportional to pulse duration (duh!)
- Can normalize energy dividing by \sqrt{T} . But we rather not.

Shifted pulses (1 of 2)



- To shift a pulse we modify the argument $\Rightarrow \Box_T(t-\tau)$
 - \Rightarrow The pulse is now centered at τ (t = τ is as t = 0 before)



• Inner product of two pulses with disjoint support ($\tau > T$)

$$\langle \sqcap_{\mathcal{T}}(t), \sqcap_{\mathcal{T}}(t- au)
angle := \int_{-\infty}^{\infty} \sqcap_{\mathcal{T}}(t) \sqcap_{\mathcal{T}}(t- au) = 0$$

> The signals are orthogonal, and indeed, "unrelated" to each other

Shifted pulses (2 of 2)



• Inner product of two pulses with overlapping support ($\tau > T$)

$$\left\langle \sqcap_{\mathcal{T}}(t), \sqcap_{\mathcal{T}}(t- au)
ight
angle := \int_{-\infty}^{\infty} \sqcap_{\mathcal{T}}(t) \sqcap_{\mathcal{T}} (t- au)$$

• The signals overlap between $\tau - T/2$ and T/2. Thus

$$\langle \Box_{\tau}(t), \Box_{\tau}(t-\tau) \rangle = \int_{\tau-\tau/2}^{\tau/2} (1)(1) dt = \frac{T}{2} - \left(\tau - \frac{T}{2}\right) = T - \tau$$



Inner product is proportional to the relative overlap

 \Rightarrow which is, indeed, how much the signals are "related" to each other



- \blacktriangleright Inner product and energy are indefinite integrals $\ \Rightarrow$ need not exist
- Complex exponential of frequency f is e_f with $e_f(t) = e^{j2\pi ft}$
- Since they have unit modulus $(|e_f(t)| = |e^{j2\pi ft}| = 1)$, their energy is

$$\|oldsymbol{e}_{oldsymbol{f}}\|^2 := \int_{-\infty}^\infty |oldsymbol{e}_{oldsymbol{f}}(t)|^2 dt = \int_{-\infty}^\infty 1 dt = \infty$$

Inner product of complex exponentials not defined ("keeps oscillating")

$$\langle e_{f}, e_{g} \rangle := \int_{-\infty}^{\infty} e_{f}(t) e_{g}^{*}(t) dt = \int_{-\infty}^{\infty} e^{j2\pi ft} e^{-j2\pi gt} dt = \int_{-\infty}^{\infty} e^{j2\pi (f-g)t} dt \Rightarrow \nexists$$

This is a problem because we can't talk about orthogonality
 Still, a complex exponential is much more like itself than another



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▶ The Fourier transform of x is the function $X : \mathbb{R} \to \mathbb{C}$ with values

$$X(f) := \int_{-\infty}^{\infty} x(t) e^{-j2\pi ft} dt$$

- We write $X = \mathcal{F}(x)$. All values of X depend on all values of x
- ▶ Integral need not exist ⇒ Not all signals have a Fourier transform
- The argument f of the Fourier transform is referred to as frequency
- Or, define e_f with values $e_f(t) = e^{j2\pi f t}$ to write as inner product

$$X(f) = \langle x, e_f \rangle = \int_{-\infty}^{\infty} x(t) e_f^*(t) dt$$

▶ Both, time and frequency are real ⇒ domain is infinite and dense ⇒ This is an analytical tool, not a computational tool (as the DFT)



▶ Since pulse is not null only when $T/2 \le t \le T/2$ we reduce X(f) to

$$X(f) := \int_{-\infty}^{\infty} \Box_T(t) e^{-j2\pi ft} dt = \int_{-T/2}^{T/2} e^{-j2\pi ft} dt$$

For $f \neq 0$, the primitive of $e^{-j2\pi ft}$ is $(-1/j2\pi f)e^{-j2\pi ft}$, which yields

$$X(f) = \left[\frac{-e^{-j2\pi f T/2}}{j2\pi f} - \frac{-e^{+j2\pi f T/2}}{j2\pi f}\right] = \frac{\sin(\pi f T)}{\pi f}$$

• Where we used $e^{j\pi fT} - e^{-j\pi fT} = 2j\sin(\pi fT)$

For
$$f = 0$$
 we have $e^{-j2\pi ft} = 1$ and $X(f)$ reduces to $\Rightarrow X(f) = T$



► Transform is important enough to justify definition of sinc function

$$\frac{\operatorname{sinc}(u) = \frac{\operatorname{sin}(u)}{u}}{\operatorname{sinc}(u) = 1} \quad \text{for } u \neq 0$$

- Value at origin, sinc(0) = 1, makes the function continuous
- With this definition and $f \neq 0$ we can write the pulse transform as

$$X(f) = \frac{\sin(\pi f T)}{\pi f} = T \frac{\sin(\pi f T)}{\pi f T} = T \operatorname{sinc}(\pi f T)$$

• Which is also true for f = 0 because $X(0) = T \operatorname{sinc}(\pi 0T) = T$



Fourier transform of pulse of width T is sinc with null crossings $\frac{k}{T}$



• Most of the Fourier Transform energy is between -1/T and 1/T

$$\int_{-1/T}^{1/T} |X(f)|^2 df = \int_{-1/T}^{1/T} |T\operatorname{sinc}(\pi fT)|^2 df \approx 0.90 T = 0.90 || \Box_T (t) ||^2$$

• Transforms of wider pulses are more concentrated around f = 0



Consistent with interpretation that shorter pulses are faster varying



• Transforms of wider pulses are more concentrated around f = 0



Consistent with interpretation that shorter pulses are faster varying



• Transforms of wider pulses are more concentrated around f = 0



Consistent with interpretation that shorter pulses are faster varying





- Let's compute a Fourier transform by approximating the integral
- Use samples spaced by T_s time units

$$X(f) = \int_{-\infty}^{\infty} x(t) e^{-j2\pi f t} dt \approx T_s \sum_{-\infty}^{\infty} x(nT_s) e^{-j2\pi f nT_s}$$

▶ Still not computable \Rightarrow consider only *N* samples from 0 to *N* − 1

$$X(f) \approx T_s \sum_{k=0}^{N-1} x(nT_s) e^{-j2\pi f nT_s}$$

▶ This is true for all frequencies. Consider frequencies $f = (k/N)f_s$

$$X\left(\frac{k}{N}f_{s}\right) \approx T_{s}\sum_{k=0}^{N-1} x(nT_{s})e^{-j2\pi(k/N)f_{s}nT_{s}} = T_{s}\sum_{k=0}^{N-1} x(nT_{s})e^{-j2\pi kn/N}$$

Definition of the DFT of a discrete signal (up to constants)

DFT as approximation of Fourier transform



• Define \tilde{x} with $\tilde{x}(n) = x(nT_s)$. The DFT of $\tilde{X} = \mathcal{F}(\tilde{x})$ has components

$$\tilde{X}(k) = \frac{1}{\sqrt{N}} \sum_{k=0}^{N-1} \tilde{x}(n) e^{-j2\pi kn/N} = \frac{1}{\sqrt{N}} \sum_{k=0}^{N-1} x(nT_s) e^{-j2\pi kn/N} = \frac{1}{T_s \sqrt{N}} X\left(\frac{k}{N} f_s\right)$$



- Can then aproximate Fourier transform as $\Rightarrow X\left(\frac{k}{N}f_s\right) \approx T_s \sqrt{N}\tilde{X}(k)$
- Approximation becomes equality at infinity and beyond $(N \to \infty, T_s \to 0)$



- Complex exponential of frequency $f_0 \Rightarrow e_{f_0}(t) = e^{j2\pi f_0 t}$
- Use inner product form to write the components of $X = \mathcal{F}(e_{f_0})$ as

$$X(f) = \langle x, e_f \rangle = \langle e_{f_0}, e_f \rangle$$

- ▶ We've seen that $\langle e_{f_0}, e_f \rangle = \infty$ if $f = f_0$ and oscillates (∄) if $f \neq f_0$
- The complex exponential does not have a Fourier transform
 Happens because energy of complex exponentials is not finite
- ▶ Truncate to $T/2 \le t \le T/2 \implies$ multiply by square pulse $\sqcap_T(t)$

$$\widetilde{e}_{f_0T}(t) := e_{f_0}(t) \sqcap_T (t) = e^{j2\pi f_0 t} \sqcap_T (t)$$



- ▶ Truncated exponential not null only when $T/2 \le t \le T/2$ (pulse)
- ▶ Then, the Fourier transform $\tilde{X}_T(f) := \mathcal{F}(\tilde{e}_{f_0T})$ is given by

$$\tilde{X}(f) := \int_{-\infty}^{\infty} e^{j2\pi f_0 t} \sqcap_{T} (t) e^{-j2\pi f t} dt = \int_{-T/2}^{T/2} e^{j2\pi f_0 t} e^{-j2\pi f t} dt = \int_{-T/2}^{T/2} e^{-j2\pi (f-f_0)t} dt$$

Same as pulse transform, except for frequency shift in exponent

For $f \neq f_0$, primitive of $e^{-j2\pi ft}$ is $(-1/j2\pi(f - f_0))e^{-j2\pi(f - f_0)t}$. Thus

$$\tilde{X}(f) = \left[\frac{-e^{-j2\pi(f-f_0)T/2}}{j2\pi(f-f_0)} - \frac{-e^{+j2\pi(f-f_0)T/2}}{j2\pi(f-f_0)}\right] = \frac{\sin(\pi(f-f_0)T)}{\pi(f-f_0)}$$

► For $f = f_0$ we have $e^{-j2\pi(f-f_0)t} = 1$ and $\tilde{X}(f)$ reduces to $\Rightarrow \tilde{X}(f) = T$





Fourier transform of truncated complex exponential is shifted sinc

 $\tilde{X}(f) = T \operatorname{sinc}(\pi(f - f_0)T)$

Transform, (centered at frequency $f_0 = 1$)



- As $T \to \infty$ truncated exponential approaches exponential
 - \Rightarrow And shifted sinc becomes infinitely tall \Rightarrow delta function





Fourier transform of truncated complex exponential is shifted sinc

 $\tilde{X}(f) = T \operatorname{sinc}(\pi(f - f_0)T)$

Transform, (centered at frequency $f_0 = 1$)



- \blacktriangleright As $\mathcal{T} \rightarrow \infty$ truncated exponential approaches exponential
 - \Rightarrow And shifted sinc becomes infinitely tall \Rightarrow delta function





▶ Fourier transform of truncated complex exponential is shifted sinc

 $\tilde{X}(f) = T \operatorname{sinc}(\pi(f - f_0)T)$ Transform, (centered at frequency $f_0 = 1$)

2.0 1.5 1.0 0.5 0 -1 0 1 2 3 frequency f (in hertz)

- As $T \to \infty$ truncated exponential approaches exponential
 - \Rightarrow And shifted sinc becomes infinitely tall \Rightarrow delta function



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• Given a transform X, the inverse Fourier transform is defined as

$$x(t) := \int_{-\infty}^{\infty} X(f) e^{j2\pi f t} df$$

- We denote the inverse transform as $x = \mathcal{F}^{-1}(X)$
- Sign in the exponent changes with respect to Fourier transform
- Can write as inner product $\Rightarrow x(t) = \langle X, e_{-t} \rangle \ (e_{-t}(f) = e^{-j2\pi ft})$
- ▶ As in the case of the iDFT, this is not the most useful interpretation



Theorem

The inverse Fourier transform \tilde{x} of the Fourier transform X of a given signal x is the given signal x

$$\tilde{\mathbf{x}} = \mathcal{F}^{-1}(\mathbf{X}) = \mathcal{F}^{-1}[\mathcal{F}(\mathbf{x})] = \mathbf{x}$$

Signals with Fourier transforms can be written as sums of oscillations

$$x(t) = \int_{-\infty}^{\infty} X(f) e^{j2\pi ft} df \approx (\Delta f) \sum_{n=\infty}^{\infty} X(f_n) e^{j2\pi f_n t}$$

This is conceptual, not literal (as was the case in discrete signals)



• X(f) determines the density of frequency f in the signal x(t)

$$x(t) pprox \sum_{n=\infty}^{\infty} (\Delta f) X(f_n) e^{j 2 \pi f_n t}$$

It represents relative contribution (as opposed to absolute)



- Signal on left accumulates mass at low frequencies (changes slowly)
- Signal on right accumulates mass at high frequencies (changes fast)



• We want to show $\Rightarrow \tilde{x} = \mathcal{F}^{-1}(X) = \mathcal{F}^{-1}[\mathcal{F}(x)] = x$. Use definitions

From definition of inverse transform of $X \Rightarrow \tilde{x}(\tilde{t}) := \int_{-\infty}^{\infty} X(f) e^{j2\pi f\tilde{t}} df$

From definition of transform of $x \Rightarrow X(f) := \int_{-\infty}^{\infty} x(t)e^{-j2\pi ft} dt$

Substituting expression for X(f) into expression for $\tilde{x}(\tilde{t})$ yields

$$\tilde{x}(\tilde{t}) = \int_{-\infty}^{\infty} \left[\int_{-\infty}^{\infty} x(t) e^{-j2\pi f t} dt \right] e^{j2\pi f \tilde{t}} df$$

Repeating steps done for DFT and iDFT with integrals instead of sums



$$\tilde{x}(\tilde{t}) = \int_{-\infty}^{\infty} x(t) \left[\int_{-\infty}^{\infty} e^{j2\pi f\tilde{t}} e^{-j2\pi ft} df \right] dt$$

- Pulled x(t) out because it doesn't depend on k
- Innermost integral is the inner product between $e_{\tilde{t}}$ and e_t .

$$\int_{-\infty}^{\infty} e^{j2\pi f\tilde{t}} e^{-j2\pi ft} df = \langle e_{\tilde{t}}, e_t \rangle$$

- Up until now we repeated same steps we did for DFT and iDFT
- But we encounter a problem $\Rightarrow \langle e_{\tilde{t}}, e_t \rangle$ does not exist (infinity, oscillates)
- ► To exchange integration order, all integrals have to exist. But one doesn't
 - \Rightarrow It is mathematically incorrect to interchange the order of integration





Replace infinite summation boundaries with finite summation boundaries

$$\tilde{x}(\tilde{t}) \stackrel{\mathsf{F}\to\infty}{=} \int_{-\infty}^{\infty} x(t) \left[\int_{-F/2}^{F/2} e^{j2\pi f\tilde{t}} e^{-j2\pi ft} df \right] dt$$

- Eventually, we need to take $F \to \infty$, but not yet.
- All integrals exist now. Innermost one is a sinc (truncated exponential)

$$\int_{-F/2}^{F/2} e^{j2\pi f\tilde{t}} e^{-j2\pi f\tilde{t}} df = F \operatorname{sinc}(\pi(t-\tilde{t})F)$$

Substitute sinc for innermost integral on previous expression

$$\tilde{x}(\tilde{t}) \stackrel{\text{F} \to \infty}{=} \int_{-\infty}^{\infty} x(t) \bigg[F \operatorname{sinc}(\pi(t-\tilde{t})F) \bigg] dt$$



- ► take the limit formally $\Rightarrow \tilde{x}(\tilde{t}) = \lim_{F \to \infty} \int_{-\infty}^{\infty} x(t) \left[F \operatorname{sinc}(\pi(t \tilde{t})F) \right] dt$
- The sinc function is centered at time $t = \tilde{t}$
- ▶ The sinc becomes infinitely tall and thin as we take $F \to \infty$
- Can then take $x(\tilde{t})$ outside of the integral (only "meaningful" value)

$$\tilde{x}(\tilde{t}) = \lim_{F \to \infty} x(\tilde{t}) \int_{-\infty}^{\infty} F \operatorname{sinc}(\pi(t - \tilde{t})F) dt$$

- The sinc function has unit integral $\Rightarrow \int_{-\infty}^{\infty} F \operatorname{sinc}(\pi(t-\tilde{t})F) = 1$
- We then have $\tilde{x}(\tilde{t}) = x(\tilde{t})$ and $\tilde{x} = x$ as we wanted to show

Fourier transform pairs

- ▶ Symmetry between transform and inverse \Rightarrow Transform pairs
- ▶ Interpret given function z as signal. Fourier transform $X = \mathcal{F}(z)$ is

$$X(f) = \int_{-\infty}^{\infty} z(t) e^{-j2\pi f t} dt$$

• Conjugate z and interpet z^* as a transform. Inverse $x = \mathcal{F}^{-1}(z^*)$ is

$$x(t) = \int_{-\infty}^{\infty} z^*(f) e^{j2\pi ft} df = \left[\int_{-\infty}^{\infty} z(f) e^{-j2\pi ft} df\right]^{2}$$

Same integrals except for switch of integration index and argument

 $X(f) = x^*(t)$, when f = t

• X is transform of z and z is transform of $X^* \equiv x^* \Rightarrow$ They are a pair

 \Rightarrow Conjugation unnecessary when signal and transform are real




Square of length $T \Rightarrow$ Sinc with zero crossings at k/T, $Tsinc(\pi fT)$



▶ Sinc with zero crossings at k/F, $Tsinc(\pi Ft) \Rightarrow$ Square of length F



Transform of sinc pulse is difficult to compute through direct operation



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$$\delta(t) := \lim_{F \to \infty} Fsinc(\pi Ft)$$

• Limit is
$$\delta(t) = \infty$$
 for $t = 0$





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• Limit is
$$\delta(t) = \infty$$
 for $t = 0$



Sequence of progressively taller sinc pulses



- Define the continuous time delta function as the limit of a sinc pulse
- $\delta(t) := \lim_{F \to \infty} Fsinc(\pi Ft)$ 4F 3F • Limit is $\delta(t) = \infty$ for t = 02F But does not exist for other t \Rightarrow Oscillates between $\pm 1/\pi t$



- On second thought, maybe we should use a different definition
- Intuitively, we want to say that the delta function is

 \Rightarrow Infinity for $t = 0 \Rightarrow \delta(t) = \infty$ for t = 0

- \Rightarrow Null for all other $t \Rightarrow \delta(t) = 0$ for $t \neq 0$
- But the question is what can we say mathematically? \Rightarrow Integrate

- ► Integrate the product of a signal with a sinc that is thin and tall ⇒ Recovers the value of the signal at time t = 0
- Since x(0) multiplies most of sinc mass
 $\int_{-\infty}^{\infty} x(t)Fsinc(\pi Ft)dt \approx x(0)$ Can write formally as
 $\lim_{F \to \infty} \int_{-\infty}^{\infty} x(t)Fsinc(\pi Ft)dt = x(0)$
- Observe that integral is the inner product of x with sinc. Then

$$\lim_{F\to\infty} \langle x, Fsinc(\pi Ft) \rangle = x(0)$$

Inner product of a signal with arbitrarily tall sinc is its value at zero





 \blacktriangleright Define delta function as the entity δ that has this property. I.e., if

 $\langle x, \delta \rangle = x(0)$

- for any signal x, we say that δ is a delta function
- In terms of integrals we write $\Rightarrow \int_{-\infty}^{\infty} x(t)\delta(t)dt = x(0)$
- Is the delta function a function? \Rightarrow Of course not
- \blacktriangleright We say that δ is a distribution or generalized function
- Abstract entity without meaning until we pass through an integral
 ⇒ Can't observe directly, but can observe its effect on other signals
- Can define orthogonality and transforms of complex exponentials



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- Consider complex exponentials of frequencies f and g
 - \Rightarrow Frequency $f \Rightarrow e_f(t) = e^{j2\pi ft}$. Frequency $g \Rightarrow e_g(t) = e^{j2\pi gt}$
- We define their inner product $\langle e_f, e_g \rangle$ as the delta function $\delta(f g)$

$$\langle e_f, e_g \rangle = \delta(f - g)$$

- ► This is a definition, not a derivation. We are accepting it to be true.
- ► If it is a definition: Does it make sense? What's its meaning?



- Complex exponentials don't have a mutual inner product.
- ▶ But truncated exponentials $e_{f,T}$ and e_{gT} do have a mutual product ⇒ Multiply by \sqcap_T . Make signal null for t > T/2 and t < T/2
- Can write inner product of truncated signals as

$$\langle e_{fT}, e_{gT} \rangle := \int_{-T/2}^{T/2} e_f(t) e_g^*(t) dt = \int_{-T/2}^{T/2} e^{j2\pi ft} e^{-j2\pi gt} dt = \int_{-T/2}^{T/2} e^{j2\pi (f-g)t} dt$$

• Integral above resolves to a sinc with zero crossings at k/T

$$\langle e_{fT}, e_{gT} \rangle = T \operatorname{sinc} [\pi (f - g) T]$$

- \blacktriangleright As $\mathcal{T} \rightarrow \infty$ truncated signals approach non-truncated counterparts...
- ...and the sinc limit is our first attempt at defining $\delta(f-g)$
- Definition didn't work. But we are looking for sense, not meaning



- Delta function is not observable directly, only after integration
- For an arbitrary given signal X(f) we must have

$$\int_{-\infty}^{\infty} X(f) \langle e_{fT}, e_{gT} \rangle df = \int_{-\infty}^{\infty} X(f) \delta(f-g) df = X(g)$$

Equivalently, we can write in terms of integrals

$$\int_{-\infty}^{\infty} \int_{-\infty}^{\infty} X(f) e^{j2\pi ft} e^{-j2\pi gt} dt df = X(g)$$

OK, fine, but really, stop messing and tell us what it means

 $\Rightarrow \text{ When } f = g \ \Rightarrow \langle e_f, e_f \rangle = \infty. \text{ When } f \neq g \ \Rightarrow \langle e_f, e_g \rangle = 0$

Can use for intuitive reasoning, but not for mathematical derivations



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- Again, we can define, not derive, the Fourier transform of e_g
- Denote as $X_g := \mathcal{F}(e_g)$ the transform of e_g . We define X_g as

 $X_g(f) = \delta(f-g)$



We draw delta functions with an arrow pointing to the sky



- Does it make sense to have $X_g(f) = \delta(f g)$
- ► Yes ⇒ Transform definition consistent with orthogonality definition

$$X_g(f) = \langle e_g, e_f \rangle = \delta(f - g)$$

 \blacktriangleright Yes \Rightarrow Definition is consistent with definition of inverse transform

$$e_g(t) = \int_{-\infty}^{\infty} X_g(f) e^{j2\pi ft} df = \int_{-\infty}^{\infty} \delta(f-g) e^{j2\pi ft} df = e^{j2\pi gt}$$

- Making $X_g(f) = \delta(f g)$ maintains Fourier analysis coherence
- Definition has clear, albeit, disappointingly trivial meaning
- ▶ Exponential of freq. g can be written as exponential of freq. g



- Denote as X_u the transform of the shifted delta function $\delta(t-u)$
- This one we can compute \Rightarrow Complex exponential of frequency u

$$X_u(f) = \int_{-\infty}^{\infty} \delta(t-u) e^{-j2\pi ft} dt = e^{-j2\pi fu} = e_{-u}(f)$$



It is the inverse we need to define as a delta function centered at u

The delta - constant transform pair



- When frequencies are null we have constants and unshifted deltas
- ► Transform of $x(t) = \delta(t) \Rightarrow X(f) = 1$. Transform of $x(t) = 1 \Rightarrow X(f) = \delta(f)$





► To find Fourier transform of cosine write as difference of exponentials

$$\cos(2\pi gt) = \frac{1}{2} \Big[e^{j2\pi gt} + e^{-j2\pi gt} \Big]$$

Since Fourier is a linear operator we transform each of the summands

$$X(f) = \frac{1}{2} \Big[\delta(f-g) + \delta(f+g) \Big]$$



▶ Pair of deltas of "height 1/2" at (opposite) frequencies $\pm g$



Continuous time signals

Fourier transform

Inverse Fourier transform

Delta function

Generalized orthogonality

Generalized Fourier transforms

Properties of the Fourier transform

Convolution



- ► Fourier transform is conjugate symmetric, linear, and conserves energy
- Transforms of real signals satisfy $\Rightarrow X(-k) = X^*(k)$
- Linearity $\Rightarrow \mathcal{F}(ax + by) = a\mathcal{F}(x) + b\mathcal{F}(y)$

• Energy
$$\Rightarrow \int_{-\infty}^{\infty} |x(t)|^2 dt = ||x||^2 = ||X||^2 = \int_{-\infty}^{\infty} |X(t)|^2 dt$$

- ▶ Not surprising, Fourier transform and DFT are conceptually identical
- Properties follow from properties of inner products and orthogonality
- Both transforms are projections on complex exponentials (inner product)
- And both project onto sets of orthogonal signals



Theorem

The Fourier transform $X = \mathcal{F}(x)$ of a real signal x is conjugate symmetric

 $X(-f) = X^*(f)$

- ► For real signals only positive half of spectrum carries information
- Conjugate symmetry implies that X(-f) and $X^*(f)$ are such that...
 - \Rightarrow Real parts are equal \Rightarrow Re(X(f)) = Re(X(-f))
 - \Rightarrow Imaginary parts are opposites \Rightarrow Im (X(f)) = Im (X(-f))
 - \Rightarrow Moduli are equal \Rightarrow |X(f)| = |X(-f)|



Proof.

• Write the Fourier transform X(-k) using its definition

$$X(-f) := \int_{-\infty}^{\infty} x(t) e^{-j2\pi(-f)t} dt$$

- When the signal is real, its conjugate is itself $\Rightarrow x(n) = x^*(n)$
- Conjugating a complex exponential \Rightarrow changing the exponent's sign

• Can then rewrite
$$\Rightarrow X(-f) := \int_{-\infty}^{\infty} x^*(t) \left(e^{-j2\pi ft}\right)^* dt$$

Integration and multiplication can change order with conjugation

$$X(-f) = \left[\int_{-\infty}^{\infty} x^*(t) \left(e^{-j2\pi ft}\right)^* dt\right]^* = X^*(f)$$



Theorem

The Fourier transform of a linear combination of signals is the linear combination of the respective Fourier transforms of the individual signals,

 $\mathcal{F}(ax + by) = a\mathcal{F}(x) + b\mathcal{F}(y).$

Proof.

• Let $Z := \mathcal{F}(ax + by)$. From the Fourier transform definition

$$Z(f) = \int_{-\infty}^{\infty} \left[a x(t) + b y(t) \right] e^{-j2\pi f t} dt$$

Expand the product, reorder terms, identify transforms of x and y

$$Z(f) = a \int_{-\infty}^{\infty} x(t) e^{-j2\pi ft} dt + b \int_{-\infty}^{\infty} y(t) e^{-j2\pi ft} dt = a X(f) + b Y(f) \quad \Box$$



Theorem (Parseval)

Let $X = \mathcal{F}(x)$ be the Fourier transform of signal x. The energies of x and X are the same, i.e.,

$$\int_{-\infty}^{\infty} |x(t)|^2 dt = ||x||^2 = ||X||^2 = \int_{-\infty}^{\infty} |X(f)|^2 df$$

It follows that X(f) is the energy density concentrated around f
 E.g., removing frequency component ≡ remove corresponding energy

We omit proof as it is analogous to DFT case. Need to use finite integration region and take limit after exchanging order of integration. Not worth repeating.



- Two more properties we didn't study for DFTs
 They (sort of) hold for DFTs, but are difficult to explain
- Time shift \Rightarrow multiplication by complex exponential in frequency
- Multiplication by complex exponential in time \Rightarrow Shift in frequency
- ▶ Properties are dual of each other ⇒ inverse transform symmetry
 ⇒ If one holds the other has to be true



- Given signal x and shift τ define shifted signal $x_{\tau} \Rightarrow x_{\tau} = x(t \tau)$
- Fourier transform of x is $X = \mathcal{F}(x)$. Transform of x_{τ} is $X_{\tau} = \mathcal{F}(x_{\tau})$.

Theorem

A time shift of τ units in the time domain is equivalent to multiplication by a complex exponential of frequency $-\tau$ in the frequency domain

$$x_{\tau} = x(t-\tau) \qquad \Longleftrightarrow \qquad X_{\tau}(f) = e^{-j2\pi f \tau} X(f)$$

► The phase of X(f) changes, but the modulus remains the same $|X_{\tau}(f)| = |e^{-j2\pi f\tau}X(f)| = |e^{-j2\pi f\tau}| \times |X(f)| = |X(f)|$

► Useful in signal detection ⇒ Don't have to compare different shifts



Proof.

- Shifted signal transform $\Rightarrow X_{\tau}(f) = \int_{-\infty}^{\infty} x(t-\tau)e^{-j2\pi ft}dt$
- Change of variables $u = t \tau$. Separate exponent in two factors

$$X_{\tau}(f) = \int_{-\infty}^{\infty} x(u) e^{-j2\pi f(u+\tau)} du = \int_{-\infty}^{\infty} x(u) e^{-j2\pi f\tau} e^{-j2\pi fu} du$$

• Pull the term $e^{-j2\pi f\tau}$ out of the integral. Identify X(f)

$$X_{\tau}(f) = e^{-j2\pi f\tau} \int_{-\infty}^{\infty} x(u) e^{-j2\pi fu} du = e^{-j2\pi f\tau} X(f)$$



- ► For signal x and freq. g define modulated signal $\Rightarrow x_g = e^{-j2\pi gt} x(t)$
- Fourier transform of x is $X = \mathcal{F}(x)$. Transform of x_g is $X_{\tau} = \mathcal{F}(x_g)$.

Theorem

A multiplication by a complex exponential of frequency g in the time domain is equivalent to a shift of g units in the frequency domain

$$x_g = e^{j2\pi gt} x(t) \qquad \Longleftrightarrow \qquad X_g(f) = X(f-g)$$

- Dual of time shift result \Rightarrow Proof not really necessary
- Principle behind transmission of signals on electromagnetic spectrum



- ▶ Signal x has bandwidth $W \Rightarrow X(f) = 0$ for $f \notin [-W/2, W/2]$
- ► Multiplying by complex exponential shifts spectrum to the right ⇒ Re-center spectrum at frequency g



• Can recover signal x by multiplying with conjugate frequency $e^{-j2\pi gt}$

Modulation of multiple bandlimited signals



► Modulate two signals with bandwidth W using frequencies g₁ and g₂ ⇒ Spectrum of x recentered at g₁. Spectrum of y recentered at g₂



- Sum up to construct signal $z(t) = x_{g_1}(t) + y_{g_2}(t)$
 - \Rightarrow Can we recover x and y from mixed signal z? \Rightarrow Yes



▶ No spectral mixing if modulating frequencies satisfy $g_2 - g_1 > W$



- To recover x multiply by conjugate frequency $e^{-j2\pi g_1 t}$
- ▶ And eliminated all frequencies outside the interval [-W/2, W/2]
- To recover y multiply by conjugate frequency $e^{-j2\pi g_2 t}$
- ▶ And eliminated all frequencies outside the interval [-W/2, W/2]



Continuous time signals

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Properties of the Fourier transform

Convolution



- Both, Fourier transforms and DFTs are:
 - \Rightarrow Conjugate symmetric, linear, & conserve energy
- The Fourier transform also satisfies shift and modulation theorems
 They also (sort of) hold for DFTs (although we haven't shown)
 As they should, DFTs are close to Fourier transforms
- A sixth property of Fourier transforms, also sort of true for DFTs
 ⇒ Convolution in time equivalent to multiplication in frequency

Convolution



- Given signal x with values x(t) and signal h with values h(t)
- Convolution of x with h is the signal y = x * h with values

$$[x * h](t) = y(t) = \int_{-\infty}^{\infty} x(u)h(t-u) \, du$$

• Operation is commutative $\Rightarrow [x * h] \equiv [h * x]$

$$[h*x](t) = \int_{-\infty}^{\infty} h(u) \times (t-u) \, du = \int_{-\infty}^{\infty} h(t-v) \times (v) \, dv = [x*h](t)$$

• Still, prefer to interpret roles of x and h as asymmetric \Rightarrow x hits h

$$x \longrightarrow h$$
 $y = x * h$
Convolution with delta functions



- Convolution with $x(t) = \delta(t) \Rightarrow y(t) = \int_{-\infty}^{\infty} \delta(u)h(t-u) du = h(t)$
- ▶ Hitting *h* with delta function produces convolution output $y \equiv h$



• Convolution with delayed delta $x(t) = \delta(t - s)$ (u = s in integrand)

$$\mathbf{y}(t) = \int_{-\infty}^{\infty} \delta(\mathbf{u} - \mathbf{s}) h(t - u) \, d\mathbf{u} = h(t - s)$$

Hitting h with delayed delta produces delayed h as output



• Convolution with scaled delta function $x(t) = \alpha \delta(t)$

$$\mathbf{y}(t) = \int_{-\infty}^{\infty} \alpha \delta(u) h(t-u) \, du = \alpha \int_{-\infty}^{\infty} \delta(u) h(t-u) \, du = \alpha h(t)$$

• Convolution with scaled and delayed delta $x(t) = \alpha \delta(t-s)$

$$\mathbf{y}(t) = \int_{-\infty}^{\infty} \alpha \delta(\mathbf{u} - \mathbf{s}) h(t - u) \, d\mathbf{u} = \alpha \int_{-\infty}^{\infty} \delta(\mathbf{u} - \mathbf{s}) h(t - u) \, d\mathbf{u} = \alpha h(t - s)$$



Convolution with scaled and delayed delta is scaled and delayed h



• Approximate convolution with Riemann sum (sampling at $u = u_n$)

$$y(t) = \int_{-\infty}^{\infty} x(u)h(t-u) \, du \approx T_s \sum_{n=-\infty}^{\infty} x(u_n)h(t-u_n)$$

► For each $u_n \Rightarrow$ Scale h(t) by $x(u_n)$ to produce $x(u_n)h(t)$ \Rightarrow Shift to time u_n to produce $x(u_n)h(t - u_n)$

Sum over all possible $u_n \Rightarrow$ integrate over all u, in the limit





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▶ Sum over all possible $u_n \Rightarrow$ integrate over all u, in the limit





Theorem (Convolution theorem)

Given signals x and y with transforms $X = \mathcal{F}(x)$ and $Y = \mathcal{F}(y)$. The Fourier transform $Z = \mathcal{F}(z)$ of the convolved signal z = x * y is the product Z = XY

$$z = x * y \iff Z = XY$$

- Convolution in time domain \equiv to multiplication in frequency domain
- ▶ When we convolve signals x and y in the time domain ⇒ Their transforms are multiplied in the frequency domain
- ► When we multiply two transforms in the frequency domain ⇒ The signals get convolved in the time domain

Proof.

• Use the definition of Fourier transform to write the transform of Z as

$$Z(f) = \int_{-\infty}^{\infty} z(t) e^{-j2\pi ft} dt$$

• Use the definition of convolution to write the signal z as

$$z(t) = \int_{-\infty}^{\infty} x(u)h(t-u) \, du$$

• Substitute the expression for z(t) into expression for Z(f)

$$Y(f) = \int_{-\infty}^{\infty} \left(\int_{-\infty}^{\infty} x(u)h(t-u) \, du \right) e^{-j2\pi ft} \, dt$$





Proof.

Rewrite the nested integral as a double integral

$$Y(f) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} x(u)h(t-u)e^{-j2\pi ft} \, du \, dt$$

• Make the change of variables v = t - u and write

$$Y(f) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} x(u)h(v)e^{-j2\pi f(u+v)} du dt$$

• Write $e^{-j2\pi f(u+v)} = e^{-j2\pi f u} e^{-j2\pi f v}$ and reorder terms to obtain

$$Y(f) = \left(\int_{-\infty}^{\infty} x(u)e^{-j2\pi f u} \, du\right) \left(\int_{-\infty}^{\infty} h(v)e^{-j2\pi f v} \, dv\right)$$

• Factors on the right are the Fourier transforms X(f) and Y(f)



- ➤ Convolution in time equivalent to multiplication in frequency ⇒ Is this useful in any way? ⇒ Certainly, few facts are more useful
- Convolution theorem implies that these two systems are equivalent



▶ The lower path for design, the upper path for implementation

The signal and the noise

- Penn Renn
- ▶ There is signal and noise, but what is signal and what is noise?
- \blacktriangleright We already know answer \Rightarrow Signal discernible in frequency domain



Original signal x(t). It moves randomly, but not that much

The signal and the noise



- There is signal and noise, but what is signal and what is noise?
- We already know answer \Rightarrow Signal discernible in frequency domain



Fourier transform X(f) of original signal

▶ Filter out all frequencies above 100Hz (and below -100Hz)



Multiply spectrum with low pass filter H(f) = ⊓_W(f) with W = 200Hz ⇒ Only frequencies between ±W/2 = ±100Hz are retained



Fourier transform Y(f) = H(f)X(f) of filtered signal

This spectral operation does separate signal from noise



Multiply spectrum with low pass filter H(f) = ⊓_W(f) with W = 200Hz ⇒ Only frequencies between ±W/2 = ±100Hz are retained



Filtered signal y(t) with y = x * h and $h = \mathcal{F}^{-1}(H) = \mathcal{F}^{-1}(\Box_W)$

This spectral operation does separate signal from noise



We can implement filtering in the frequency domain

 $\Rightarrow \mathsf{Sample} \ \Rightarrow \mathsf{DFT} \ \Rightarrow \mathsf{Multiply} \ \mathsf{by} \ H(f) = \sqcap_W(f) \ \Rightarrow \mathsf{iDFT}$



We can also implement filtering in the time domain

 \Rightarrow Inverse transform of $\sqcap_W(f)$ is $h(t) = W \operatorname{sinc}(\pi W t)$

 \Rightarrow Sample (or not) \Rightarrow Implement convolution with h(t)